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Role of oxygen vacancy in metal oxides for photocatalytic CO₂ reduction

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ABSTRACT

Photocatalytic conversion of greenhouse gas CO₂ into valuable solar fuels represents a promising technology for addressing the global energy crisis and environmental issues simultaneously. In such a technology, the development of efficient photocatalysts is a central task for pushing forward the practical application of photocatalytic CO₂ reduction. Due to their low cost, high redox capability, and environmental friendliness, metal oxide-based photocatalysts have been extensively employed for CO₂ reduction. Moreover, oxygen vacancy engineering in metal oxides has gradually emerged as a versatile approach to enhancing their photocatalytic performance for CO₂ reduction. In this article, the state-of-the-art progress in oxygen vacancy engineering, including its synthesis, characterization, and recent advancement in photocatalytic CO₂ reduction, is reviewed. In particular, the roles of oxygen vacancy in promoting the three basic steps in photocatalytic CO₂ reduction, i.e., light absorption, charge separation, and surface CO₂ conversion, are discussed in detail. The current challenges and future opportunities of engineering oxygen vacancy in metal oxide-based photocatalysts for efficient CO₂ reduction are also addressed. This review aims to inspire more creative works on the rational design of oxygen vacancy that maximizes its function, hence accelerating the discovery of high-performance photocatalysts for CO₂ reduction.

1. Introduction

Since the industrial revolution, the extensive consumption of conventional fossil fuels, including coal, oil, and natural gas, has led to a sharp rise in the concentration of atmospheric carbon dioxide (CO₂) [1–3]. As a major greenhouse gas, CO₂ has brought about a series of environmental issues, including the familiar global warming and glacier melting [4–6]. It is estimated that when the concentration of atmospheric CO₂ exceeds 600 ppm, the sea level will increase by 0.4 m [7], posing a great threat to our ecosystem. To lessen the CO₂-induced global environmental issues, it has reached a consensus that the atmospheric CO₂ concentration should peak (net-zero global CO₂ emissions) by mid-century and decline in the second half of this century [8]. In this context, there is an urgent need to develop advanced CO₂ conversion technologies for a carbon-negative future society [9–14].

Since the pioneering work by Fujishima and Hongda on solar-driven water splitting [15], the utilization of inexhaustible solar energy as the

energy input for CO₂ conversion (i.e., photocatalytic CO₂ reduction) has attracted massive interest [16-18]. Through semiconductor-based photocatalytic processes [19,20], CO₂ can be converted into valuable green solar fuels [21,22], including CO, CH₄, CH₃OH, and HCOOH. Therefore, it can be considered a strategy that "kills two birds with one stone", simultaneously eliminating atmospheric CO2 and producing sustainable fuels and chemicals. In a typical semiconductor-based photocatalytic CO2 reduction process, there are three basic steps [23]: i) harvesting of solar energy by the semiconductor photocatalyst to produce photogenerated charge carriers, ii) separation and transfer of the charge carriers, and iii) surface reactions between the charge carriers and substrate molecules (i.e., CO2 and H2O). Recently, a series of metal oxide semiconductors (e.g., TiO2, ZnO, Cu2O, BiVO4, SrTiO3) have been explored as the photocatalyst for CO₂ reduction owing to their low cost, high redox capability, environmental friendliness, and excellent chemical and thermal stability [24,25]. However, the performance of pure metal oxides is still far from satisfactory, mainly due to the limited light

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absorption, severe charge carrier recombination, and sluggish surface CO_2 conversion. To mitigate these challenges, many strategies, including morphology control, defect engineering, doping, and heterojunction construction, have been proposed to achieve efficient photocatalytic CO_2 reduction [26–28]. For instance, photosensitizers have been integrated into metal oxides to extend their light-absorption range, while metal co-catalysts are loaded onto their surface to facilitate charge carrier transfer/separation and surface CO_2 conversion.

Among the various strategies [29-32], oxygen vacancy (OV) engineering has been one of the most extensively studied topics [33-35]. OV, as a typical type of point defect [36-38], inevitably exists in metal oxide-based photocatalysts [39,40], especially when the size of the metal oxide is reduced to the nanoscale [41]. OV can significantly affect the optical and electronic properties of metal oxides [42] and thus their photocatalytic performance [43]. In earlier research, OV in the metal oxide matrix was considered detrimental to the photocatalytic process because it serves as a recombination center for charge carriers [44]. Nevertheless, with the rapid development of controllable synthesis and advanced characterization techniques for OV, it has been gradually recognized that OV can potentially influence all three basic steps of semiconductor-based photocatalytic CO₂ reduction. Furthermore, the in-depth understanding of the role of OV has provided plenty of opportunities for performance optimization by OV engineering. A typical case is the rational design of OV in the bulk and surface of metal oxide photocatalysts to achieve improved light absorption and accelerated charge carrier separation simultaneously.

In this review, we focus on clarifying the roles of OV in metal oxide-based photocatalytic CO_2 reduction and demonstrating methods to make rational use of OV for enhanced photocatalytic performance (Fig. 1). Firstly, we present the fundamental concept of metal oxide-based photocatalytic CO_2 reduction and the functions of OV. Subsequently, we discuss the approaches for controllable synthesis of OV in

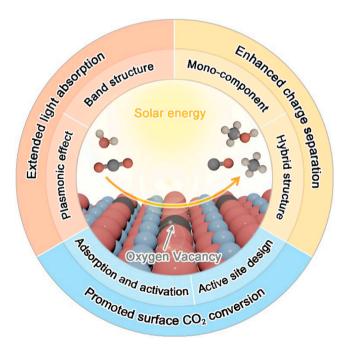


Fig. 1. Schematic illustration of the roles of OV in metal oxide-based photocatalytic CO_2 reduction. OV can extend the light absorption range of metal oxide photocatalysts by tuning their energy band structures (e.g., band positions and defect states) and inducing a plasmonic effect. The OV-induced defect state can trap photogenerated electrons, allowing enhanced charge separation in mono-component metal oxides and metal oxide-based hybrid nanostructures. Moreover, the role of OV in promoting the adsorption and activation of CO_2 provides abundant opportunities for rational active site design, benefiting the development of highly efficient photocatalysts.

metal oxides and the related characterization techniques. Furthermore, we summarize recent progress on OV engineering for enhanced photocatalytic CO_2 reduction, emphasizing the detailed roles of OV in promoting the three basic steps of semiconductor-based photocatalytic CO_2 reduction. Finally, an outlook of existing challenges and future opportunities is presented. As far as we know, this report is the first comprehensive review that clarifies the detailed roles of OV in metal oxide-based photocatalytic CO_2 reduction.

2. Fundamentals

In photocatalytic CO_2 reduction, catalytic activity and target product selectivity are two critical factors for performance evaluation. These two factors are both dependent on the three basic steps of semiconductor-based photocatalytic CO_2 reduction, light absorption, charge separation, and surface CO_2 conversion. In this sense, OV in metal oxides can significantly alter these basic steps, thus possessing immense influences on photocatalytic CO_2 reduction activity and selectivity.

2.1. Metal oxide-based photocatalytic CO2 reduction

In a typical semiconductor-based photocatalytic CO₂ reduction process (Fig. 2a), the semiconductor photocatalyst is first excited by absorbing an incident photon with energy larger than its bandgap energy [45]. Along with photon absorption, electron hopping occurs from the valence band to the conduction band of the semiconductor, leaving a hole in the valence band. The photogenerated electrons and holes are then separated and transferred from the bulk to surface sites. During this process, most of the electrons and holes are consumed by recombination (step iv in Fig. 2a). Only a fraction of them can reach the catalyst surface sites and finally be consumed by surface reactions with substrate reactants (i.e., CO2, protons, and H2O). H2O is needed to consume photogenerated holes and provide protons for the continuous conversion of CO2. In addition, given that carbon has the highest valence state in CO2 (+4), it can only accept electrons to be reduced into a variety of solar fuels, including CO (+2), HCOOH (0), CH₃OH (-2), and CH₄ (-4). In principle, the CO₂ reduction products can be controlled thermodynamically by the different redox potentials between CO2 and the various reduction products (Table 1) [17,24]. However, complications arise when surface reaction kinetics is considered. For example, although the conversion of CO2 to CH4 is the most thermodynamically favorable, the dominant product will be CO when the electron supply is insufficient due to reaction kinetic limitations. In addition, if the produced CO can be stabilized on the surface sites of the photocatalyst, it can be further combined with protons provided by H2O to produce hydrogenated products. It is worth noting that the protons can also directly accept electrons to generate hydrogen atoms, potentially contributing to H2 evolution. As such, it remains a great challenge to efficiently utilize photogenerated electrons for highly selective CO₂ reduction.

In addition to selectivity, photocatalytic activity for CO2 reduction should also be considered during photocatalyst design. To achieve efficient CO2 reduction, the semiconductor photocatalyst should have both high reduction and oxidation capabilities. Specifically, electrons on the conduction band of the semiconductor should have more negative potential than the redox potential between CO2 and the target reduction product; meanwhile, holes on the valence band should be oxidative enough to oxidize H2O into O2. In this sense, conventional widebandgap metal oxides are considered one of the most promising choices because of their suitable energy band positions (Fig. 3) [46]. However, these metal oxide photocatalysts still suffer from poor CO2 reduction activities due to the following limitations [47]. Firstly, these wide-bandgap metal oxides can only be excited by UV light (< 400 nm), which accounts for less than 5% of the solar spectrum [48]. Secondly, the charge carriers in these metal oxide photocatalysts experience severe recombination after photoexcitation. Lastly, there are limited active sites on the metal oxide surface for the adsorption and activation of CO₂,

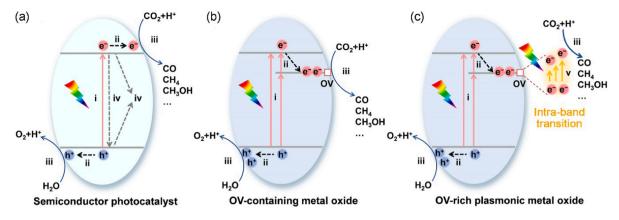


Fig. 2. Schematic illustration for (a) the overall photocatalytic process of CO_2 reduction with H_2O over a semiconductor photocatalyst and (b,c) the influences of OV on the process. The numbers i, ii, and iii represent the three basic steps in photocatalytic CO_2 reduction, i.e., light absorption, charge separation, and surface reaction. Following this, numbers iv and v represent the recombination of photogenerated electrons and holes and the intra-band transition of electrons in the OV-induced defect band, respectively. Step iv can occur either upon the generation of electrons and holes or during their migration from the bulk to surface sites. In addition, it should be predominant even in the presence of OV. However, to simplify the figure and thus facilitate reading, this step is not shown in Fig. 2b and c.

Entry	¹ Chemical equations	$^{2}E_{\rm redox}$ (V vs. NHE)		$^3 \Delta H^0$	
		$\overline{PH} = 0$	PH = 7	(kJ mol ⁻¹)	
1	$CO_2 + e^- \rightarrow \bullet CO_2^-$	- 1.44	- 1.85	/	
3	$CO_2 + 2 H^+ + 2e^- \rightarrow CO_{(g)} + H_2O_{(l)}$	-0.11	-0.52	-41.2	
2	$CO_2 + 2 H^+ + 2e^- \rightarrow HCOOH_{(1)}$	-0.20	-0.61	-31.2	
4	$CO_2 + 4 H^+ + 4e^- \rightarrow HCHO_{(g)} + H_2O_{(l)}$	0.07	- 0.48	- 121.0	
5	$CO_2 + 6 H^+ + 6e^- \rightarrow CH_3OH_{(1)} + H_2O_{(1)}$	0.03	- 0.38	- 131.0	
6	$CO_2 + 8 H^+ + 8e^- \rightarrow CH_{4(g)} + 2 H_2O_{(1)}$	0.17	-0.24	-253.0	
7	$2CO_2 + 8 H^+ + 8e^- \rightarrow CH_3COOH_{(1)} + 2 H_2O_{(1)}$	0.11	- 0.30	122.1	
8	$2CO_2 + 12 H^+ + 12e^- \rightarrow C_2H_5OH_{(I)} + 3 H_2O_{(I)}$	0.09	- 0.32	216.1	
9	$2\text{CO}_2 + 12 \text{ H}^+ + 12\text{e}^- \rightarrow \text{C}_2\text{H}_{4(g)} + 4 \text{ H}_2\text{O}_{(1)}$	0.08	- 0.33	127.8	
10	$2\text{CO}_2 + 14 \text{ H}^+ + 14\text{e}^- \rightarrow \text{C}_2\text{H}_{6(g)} + 4 \text{ H}_2\text{O}_{(1)}$	0.14	- 0.27	264.3	
11	$2 \text{ H}^+ + 2 \text{e}^- \rightarrow \text{H}_{2(g)}$	0	-0.41	0	
12	$H_2O_{(1)} + 2 h^+ \rightarrow \bullet OH_{(aq)} + 2 H^+$	2.73	2.32	71.0	
13	$2 H_2O_{(1)} + 4 h^+ \rightarrow O_{2(g)} + 4 H^+$	1.23	0.82	572.0	

¹ g: gas, l: liquid.

resulting in sluggish CO2 reduction kinetics.

2.2. OV in metal oxide photocatalysts

Fortunately, the introduction of OV into the metal oxide matrix has largely alleviated these limitations [41,43,49]. In metal oxides, OV assumes a local cave structure where an oxygen atom is removed from the oxide matrix [50]. The removal of the oxygen atom leaves two electrons at the vacancy site, and thus, the whole metal oxide system maintains a charge-neutral state (O $^{2-}$ \rightarrow O_{atom} + OV + 2 e^-). The electrons accumulated at the OV sites can be transferred to adjacent metal atoms, forming positively charged OV sites and reduced metal centers. Bao and co-workers [51] studied the influence of hydrogen treatment on the evolution of OV and Ti³⁺ sites in the TiO₂ matrix. It was found that Ti³⁺ only appeared after the treatment temperature exceeded 450 °C, while OV already formed at a lower temperature. This result corroborated that the electrons in the OV were transferred out to form Ti³⁺ when sufficient energy was supplied by the hydrogen at higher temperatures. Additionally, positively charged OV sites are formed concomitant with the electron transfer and Ti³⁺ formation. The OV in a charge state of neutral, + 1, or + 2, can be respectively denoted as V_O^x , V_O^{\bullet} , $V_O^{\bullet \bullet}$ according to the Kröger-Vink notation [50].

As an intrinsic point defect [42,52], OV commonly exists in metal oxides and can significantly alter their optical and electronic properties [53–57]. Moreover, considering that OV can be introduced in situ by light irradiation [58–60], it is expected to have a non-negligible influence on the photocatalytic process over metal oxide-based

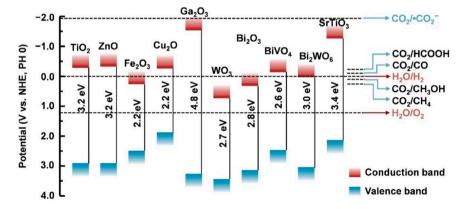


Fig. 3. Schematic illustration for energy band structures of common metal oxides. While some metal oxides (e.g., WO₃) do not have suitable conduction band positions to drive CO₂ reduction, they can form heterojunctions with other semiconductor materials and contribute to photocatalytic CO₂ conversion.

 $^{^{2}}$ $E_{\rm redox}$: redox potential

³ ΔH^0 : reaction enthalpy.

photocatalysts [37,61]. OV as an electron donor can induce a defect state (an inter-band state between conduction and valence bands) in semiconductor photocatalysts (Fig. 2b). Yim et al. used electron bombardment to vary the density of bridging OV in rutile TiO₂(110) with in situ monitoring of the bandgap state (i.e., defect state) using ultraviolet photoemission spectroscopy [62]. It was revealed that the intensity of the photoemission peak was directly proportional to the OV density, confirming the dominant contribution of OV to the defect state. The OV-induced defect state allows metal oxide to be excited by absorbing photons with energy lower than its bandgap. In addition, Zhang et al. studied the influence of OV on energy band positions of metal oxide photocatalysts using OV-rich WO₃ as an example [63]. Theoretical calculations confirmed that the valence and conduction bands of WO3 are mainly contributed by O 2p and W 5d orbitals, respectively. Accordingly, the electron transfer from OV to surrounding oxygen atoms through the coordination-unsaturated metal atoms will elevate the potential of the O 2p orbital. In other words, the OV-rich sample possesses a higher valence band maximum and thus narrower bandgap, resulting in promoted light absorption.

More interestingly, a unique electron excitation (i.e., light absorption) process has also been observed for some heavily OV-doped metal oxides (e.g., WO₃, MoO₃, and Bi₂O₃) [64–66]. In these metal oxides, the large number of electrons brought about by the heavily OV doping can absorb the incident light and experience intra-band transition in the OV-induced defect band (with a continuum of states available for ensemble electron excitation) (Fig. 2c) [67], analogous to the localized surface plasmon resonance (LSPR) excitation of free electrons in metal nanostructures (e.g., Au, Ag, and Cu) [68-70]. A key characteristic of LSPR excitation is that the optical extinction cross-section of the material will reach a maximum at certain frequencies owing to a collective excitation of electrons [68]. In other words, absorption peaks can be observed in the absorption spectra of the material, which can be monitored by UV-visible-near infrared (UV-vis-NIR) spectroscopy. Specifically, noble metal nanostructures with a high electron concentration (10²²–10²³ cm⁻³) exhibit strong LSPR peaks in the visible light region. In contrast, heavily OV-doped metal oxides typically display LSPR peaks at 700-1250 nm due to the relatively low electron concentration $(10^{21}-10^{22} \text{ cm}^{-3})$ [71]. Nonetheless, rational regulation of electron concentration in these OV-containing plasmonic metal oxides presents a promising strategy for tuning their LSPR characteristics (e.g., intensity and peak position).

In short, OV can significantly extend the light absorption range of metal oxide photocatalysts by tuning their band structures (i.e., defect state and band positions) and inducing a plasmonic effect [72]. It is worth mentioning that the OV-induced defect state can further enhance the separation of photogenerated electrons and holes by serving as an electron-trapping state [73]. Therefore, more electrons and holes can be transferred to the surface sites of the photocatalyst, contributing to CO₂ reduction and H₂O oxidation, respectively [61]. Moreover, OV can also regulate the adsorption and activation behaviors of CO₂ on the surface of metal oxide photocatalysts. Rational active site design based on OV thus enables a delicate control of CO2 reduction activity and target product selectivity. Therefore, the current studies on OV can be roughly divided into two parts: i) understanding the basics of OV including its preparation, structures and properties, and ii) rationally engineering OV for enhanced photocatalytic CO2 reduction performance. The current progress of OV engineering in photocatalytic CO2 reduction will be discussed in detail in Section 4, after a brief introduction of the synthesis and characterization techniques of OV in metal oxides.

3. Controllable synthesis and characterization of OV in metal oxides

3.1. Controllable synthesis of OV-containing metal oxides

The controllable synthesis of OV-containing metal oxides is the

prerequisite for their practical application. OV is introduced by removing an oxygen atom from the metal oxide, which can be regarded as a reduction process. Therefore, the formation of OV is determined by both the reductive environment conditions and the intrinsic reducibility of the metal oxide (i.e., OV formation energy). By extension, OV engineering in metal oxides can be approached from two aspects, the application of reductive synthesis/treatment environments and the regulation of the intrinsic OV formation energy.

3.1.1. Application of reductive environment

OV can be introduced into the metal oxide matrix during material synthesis [74]. In this case, reducing agents, such as glucose, alcohol, oxalic acid, ascorbic acid, and metallic metals, are commonly employed to provide a reductive synthesis environment for OV formation. For example, Kong et al. prepared OV-rich Bi₂WO₆ using a simple one-step ethylene glycol-assisted solvothermal method [75]. The redox reaction between ethylene glycol and Bi₂WO₆ resulted in the selective removal of oxygen atoms from the Bi₂WO₆ surface. In addition, the generated OVs can be refilled with oxygen atoms after calcinating the obtained sample in an air atmosphere at 400 °C for 4 h, as evidenced by the change in absorption in the visible light range (Fig. 4a). After material synthesis, OV can also be introduced by reduction treatment. Such OV is generally located on the surface of the material and can be accompanied by the appearance of surface disorder (amorphous layer) [76]. Huang and co-workers developed a disorder-engineered blue TiO2 sample using lithium-containing ethylenediamine (Li-EDA) as a reducing agent [77]. The obtained sample featured a unique surface disorder layer that contains numerous OVs and doped hydrogen atoms. Moreover, the thickness of the disordered layer can be controlled by simply tuning the Li dosage.

High-temperature treatment is also an effective method to create OV in metal oxides. In a high-temperature environment, the oxygen atom removal efficiency from metal oxide depends on the oxygen partial pressure. In this regard, reducing agents, such as H2, NH3, NaBH4, Mg, and Al, are commonly employed to consume oxygen and thus assist the oxygen atom removal from the metal oxide matrix. Liu et al. recently reported the synthesis of OV-rich MoO_{2-x} nanobelts by thermal treatment of MoO $_3$ nanobelts at 550 °C under a 10% H $_2$ /Ar mixture [78]. The OV-poor MnO₂ sample can be obtained by treating the OV-rich MoO_{2-x} at 200 °C in air. Similarly, Zhang and co-workers developed a universal NH₃-assisted reduction strategy to introduce OV into transition metal oxides (WO₃, MnO₂, Nb₂O₅, and MoO₃) [79]. Mechanism studies suggested that H and N atoms from NH3 were able to extract oxygen atoms in transition metal oxides to form H2O, N2, N2O, and NO (Fig. 4b). In addition, Gao et al. reported the synthesis of black TiO2 arrays using Al as the reductant [80]. Al can assist in capturing the oxygen released by metal oxides during high-temperature treatment, thus facilitating the lowering of oxygen partial pressure and formation of Ti^{3+} ions and OVs on the sample surface (Fig. 4c). The combination of solid (e.g., Mg) and gaseous (e.g., H₂) reductants can be an even more powerful technique for lowering oxygen partial pressure and introducing OV into metal oxides [81]. However, excessive OVs in the catalyst may not be beneficial for its application in photocatalytic CO2 reduction.

Despite the strong reducing capability, these externally added agents may also introduce impurities such as interstitial H and N dopants into the metal oxide matrix [82–84]. To overcome this, inert conditions (e.g., Ar and vacuum) can be used to control the oxygen partial pressure in the environment and thus effectively create OV without impurity introduction. It is worth noting that this type of treatment is generally applied to metal oxides with relatively soft metal—oxygen bonds, such as BiVO $_4$, WO $_3$, ZnO, SnO $_2$, ZnO, and CuO. High temperatures are still needed due to the weak reduction capability of the inert gases. However, too high annealing temperatures will potentially affect the morphology and phase of some metal oxides (e.g., CuO and Fe $_2$ O $_3$). On this note, Zhang and co-workers recently proposed a vacuum-assisted self-photoetching approach to synthesize OV-rich orthorhombic cobalt carbonate

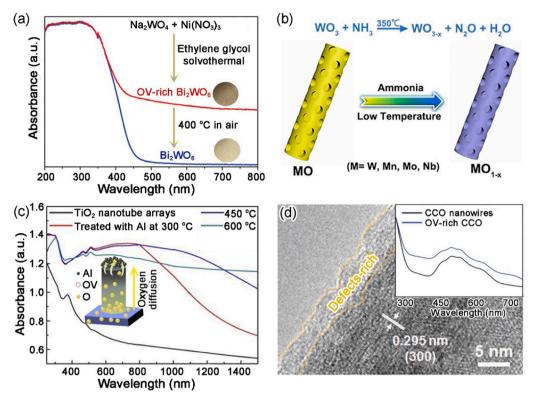


Fig. 4. (a) UV-vis diffuse reflectance spectroscopy (DRS) spectra of OV-rich Bi₂WO₆ obtained from one-step ethylene glycol-assisted solvothermal process. The inset shows the color of the samples. Reproduced with permission from Ref. [75]. Copyright 2016, The Royal Society of Chemistry. (b) Scheillustration matic for low-temperature ammonia-assisted synthetic process of OV-enriched transition metal oxides. Reproduced with permission from Ref. [79]. Copyright 2018, Elsevier B.V. (c) UV-vis DRS spectra of pristine TiO2 nanotube arrays and the TiO2 nanotube arrays treated with Al at 300, 450, and 600 °C. The inset shows the Al-assisted oxygen atom removal process in TiO2. Reproduced with permission from Ref. [80]. Copyright 2019. Elsevier B.V. High-resolution TEM image of OV-rich cobalt carbonate hydroxide nanowires (CCO, Co(CO₃)_{0.5}·0.11 H₂O) obtained vacuum-assisted photoetching approach. The inset shows UV-vis DRS spectra of pristine CCO nanowires and the OV-rich CCO nanowires. Reproduced with permission from Ref. [59]. Copyright 2021, Elsevier B.V.

hydroxide nanowires ($Co(CO_3)_{0.5}$ ·0.11 H₂O) at ambient temperature [59]. The authors demonstrated that the crack of carbonate species on the nanowire surface under light irradiation resulted in the formation of an OV-rich surface layer (i.e., coordination-unsaturated Co sites) (Fig. 4d).

In principle, many other methods that can provide sufficient energy and a reductive environment have the potential to create OV in metal oxides. For instance, plasma treatment has recently attracted exceptional attention due to its ability to create OV under mild conditions [85]. The plasma-assisted OV creation can be controlled by plasma compositions (e.g., H2, Ar, and air), power, and treatment duration. Moreover, this plasma treatment can be carried out in a fast and highly efficient way without compromising its bulk morphology. Despite its potential, the application of this technique in photocatalytic CO₂ reduction is still scarce due to the low stability of the introduced OV sites. Zhu et al. reported that a hole scavenger (triethanolamine, TEOA) was needed for CO2 reduction over an OV-rich In2O3 photocatalyst synthesized by NH3 plasma [86]. In another work, Chen et al. reported that continuous Ar plasma treatment is needed to maintain the OV on the NiO/TiO2 catalyst surface in order to achieve stable CO2 reduction [87]. Therefore, the development of novel strategies to extend the application of this time- and energy-saving technique in photocatalytic CO₂ reduction is highly desired.

3.1.2. Regulation of OV formation energy

The ease of introducing OV in a material can be described by the OV formation energy, where a lower OV formation energy means easier OV introduction and vice versa. Therefore, the regulation of OV formation energy is another effective strategy for controllable OV introduction. Shi and co-workers developed a facile Ti–OH bond-breaking approach to create OV in titania. Compared to Ti–O bonds, the energy required for the dissociation of Ti–OH bonds is much lower. Therefore, OV was efficiently introduced into titania by simply calcining $\text{TiO}_x(\text{OH})_y$ precursor in N_2 [88]. A similar method was adopted to synthesize OV-containing NiO from β -Ni(OH) $_2$ precursor [89]. In addition, Di et al. reported a strain-assisted synthesis of OV in $\text{Bi}_{12}\text{O}_{17}\text{Br}_2$ nanotubes [90].

The stretched local structure on the curved nanotube surface weakens the Bi–O bonding, and as a result, possesses a lower OV formation energy and higher OV concentration, compared to the nanoplate structure (Fig. 5a). Following this line of thought, a high-pressure torsion process is commonly applied to simultaneously introduce lattice strain and OVs in metal oxide materials [91,92]. Moreover, Huang and co-workers reported the OV formation in WO_3 film that was assisted by lattice strain arising from the thermal expansion during the photothermal reaction. The lattice strain and OV formation can collectively affect the electronic structure of the photocatalyst, resulting in enhanced performance toward photocatalytic CO_2 reduction [93].

In addition, size control of metal oxides provides the opportunity for OV engineering because OV formation is easier on the surface compared to bulk [94]. Theoretical calculations on MgO and CeO2 suggest that surface oxides possess lower OV formation energy [95,96]. A smaller size means a higher surface area and thus a higher possibility for OV formation. Based on this, Guo et al. synthesized ultra-small Bi₂Sn₂O₇ nanoparticles (~4 nm) with high surface area [97]. Resultingly, a large number of OVs was created in the ultra-small Bi₂Sn₂O₇, allowing the sample to exhibit increased light absorption in the visible light range compared to bulk Bi₂Sn₂O₇ (Fig. 5b). In addition, the same research group reported the synthesis of Bi₂MoO₆ ultrathin nanosheets (with a thickness of 1-2 nm) using layered BiOBr nanosheets as the template [98]. The results suggested the existence of abundant OVs on the material surface in the form of Bi-O vacancy pairs. Moreover, the OV confined in the ultrathin nanosheets led to a new donor level near the conduction band minimum and a higher DOS at the valence band maximum, which extended the light absorption and accelerated the charge carrier separation [99].

Recently, engineering the local structure of OV has been gradually recognized as an effective pathway for promoting OV formation. The concept of "asymmetric oxygen vacancy site" is introduced to describe an OV site with asymmetric coordination of cations $(M_1$ –OV– M_2). Under such a condition, OV sites are more easily refilled and regenerated (i.e., lower OV formation energy of the metal oxide matrix) [100]. In this regard, doping is commonly used to introduce asymmetry and thus

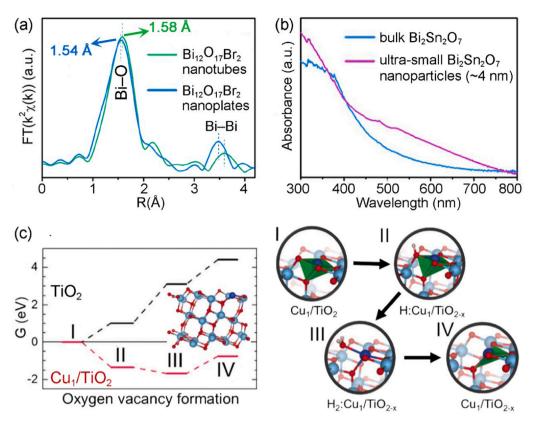


Fig. 5. (a) Fourier-transform Bi L₁-edge EXAFS spectra of Bi₁₂O₁₇Br₂ nanotubes and nanoplates, highlighting the elongated Bi-O bonds on a curved surface (i. e., with lattice strain). Reproduced with permission from Ref. [90]. Copyright 2020, American Chemical Society. (b) UV-vis DRS spectra of ultra-small Bi₂Sn₂O₇ nanoparticles, in comparison to bulk Bi₂Sn₂O₇. Reproduced with permission from Ref. [97]. Copyright 2021, Elsevier B.V. (c) Simulated energy changes of OV formation on TiO2 and Cu-doped TiO₂ (Cu₁/TiO₂). In the Cu-O₄ structure model (green), Cu, O, and Ti atoms are in blue, red, and cyan colors, respectively. Reproduced with permission from Ref. [104]. Copyright 2022, Elsevier B.V.

promote OV formation [101]. For instance, Brant et al. observed a neutral Cu^{2+} –OV complex in Cu-doped rutile TiO_2 crystalline. Experimental evidence shows that the substitution of Ti^{4+} by Cu^{2+} led to the formation of OV adjacent to the Cu^{2+} site due to charge compensation [102]. First-principle results also confirmed the formation of charge-compensating OVs at sites adjacent to the metal dopant [103]. With this in mind, Lee et al. successfully prepared an OV-rich TiO_2 photocatalyst for efficient CO_2 reduction through a facile Cu-doping strategy [104]. Theoretical calculations confirmed that the electronic interaction between Cu and TiO_2 could lower the energy for OV formation in TiO_2 (Fig. 5c), leading to spontaneous OV formation near the Cu dopant. Moreover, the Cu dopant and the introduced OV sites can jointly stabilize CO_2 reduction intermediates on the TiO_2 surface, resulting in substantially enhanced performance compared to pristine TiO_2 .

Due to the charge-compensating characteristic, doping-induced OVs typically possess much higher stability than those on pristine metal oxides. Bo et al. reported that the OV in Fe-doped TiO2 was maintained after calcination at 800 $^{\circ}\text{C}$ in air, which can hardly be expected for undoped TiO2 [105]. Similarly, Wang and co-workers also observed metal doping-facilitated OV formation in atomic Ni-loaded TiO2 catalysts during an oxidative molten salt treatment process [106]. The high stability of OV makes these metal oxides potential photocatalysts for highly efficiently photocatalytic CO2 reduction. It is also worth mentioning that non-metal doping (e.g., C and N) can likewise induce the charge-compensating effect and thus the formation and stabilization of OV on the metal oxide surface [107-109]. In addition to doping, the formation of a metal/oxide interface is another effective pathway to regulate the reducibility of metal oxides [110]. The metal/oxide interface can be formed either by loading small metal nanoparticles onto the surface of metal oxide substrates [111-113] or by coating metal oxide nanoparticles/films onto metal substrates [114]. However, this method is currently mainly used to obtain metal oxide photocatalysts for CO2 reduction with H₂ [115,116]. The use of H₂O as the reactant may not be able to regenerate OV at the metal/oxide interface during CO2

conversion. Nonetheless, the integration of metal/oxide interfaces with other modification methods (e.g., doping) may serve as a promising pathway for efficient OV formation and CO₂ reduction.

3.2. Characterization of OV in metal oxides

3.2.1. Electron paramagnetic resonance and X-ray absorption spectroscopy
Identifying the state of OV in metal oxides is essential for its
controllable synthesis and subsequent application [117–119]. In addition to the optical characterization method (i.e., UV–vis DRS), many
methods have been proposed for OV characterization (Fig. 6) [120].
Among these methods, transmission electron microscope (TEM),

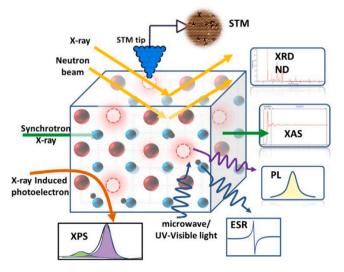


Fig. 6. Schematic illustration for the methods used to characterize OV (dashed circles) in metal oxides. Reproduced with permission from Ref. [120]. Copyright 2021, John Wiley & Sons, Inc.

electron paramagnetic resonance (EPR), X-ray absorption spectroscopy (XAS), X-ray photoemission spectroscopy (XPS), photoluminescence (PL), and positron annihilation spectrometry (PAS) are most often used. With the assistance of these characterization tools, Zhang et al. systematically investigated the OV characteristics within WO3 nanoplates [63]. In the atomic resolution high-angle annual dark-field scanning transmission electron microscopy (HAADF-STEM) image, lattice disorder and dislocation were observed in the defect-rich WO3 nanosheets (R-WO₃), indicating the formation of a defective surface (Fig. 7a). The defect types were further confirmed by EPR and W L3-edge extended X-ray absorption fine structure (EXAFS) spectra. The symmetrical EPR signal at g = 2.002 was assigned to single electron-trapped OVs on R-WO₃ (Fig. 7b). Furthermore, based on the fitting results of the W L₃-edge EXAFS spectra, the R-WO₃ displayed reduced total W-O coordination numbers compared to commercial WO3 and defect-free WO3 (D-WO₃) (Fig. 7c).

Moreover, the development of advanced in situ characterization techniques (e.g., EPR and in situ XAS) to monitor the dynamic changes of OV under light irradiation and/or during interactions with adsorbed molecules is highly desirable for understanding the role of OV [121]. For example, Xiong and co-workers employed in situ EPR to investigate the

changes of OV under light irradiation and in different atmospheres (i.e., He and N₂) [105]. The reduced EPR signal intensity suggested a decrease in the amount of single-electron trapped OV site under light irradiation in N2 (Fig. 7d). Similarly, Tang and co-workers successfully used in situ EPR to identify the role of OV as the trapping site for photogenerated electrons over Pd atoms-doped defective In₂O₃ [122]. In situ XAS is also a useful tool to trace the structure evolution of metal oxides during the catalytic reaction [123,124]. Paulus and co-workers studied the electrochemical oxidation behaviors of SrCoO_{2.5} with in situ Co K-edge XAFS spectroscopy [125]. Surprisingly, an intermediate phase (SrCoO_{2.82} ± 0.07) with oxygen ordering appeared at room temperature during the oxidation process of SrCoO_{2.5} to SrCoO₃. In another report, Park et al. [126] utilized Ce L3-edge EXAFS spectra to unravel the photo-induced reduction of Ce^{4+} to Ce^{3+} on the CeO_2 surface (Fig. 7e). The Fourier-transformed EXAFS spectra demonstrated that the CeO2 reduction is accompanied by an elongation of the Ce-O bond from 2.36 to 2.42 Å due to larger ionic radium of Ce^{3+} (1.14 Å) than Ce^{4+} (1.00 Å).

3.2.2. Positron annihilation spectroscopy

To reliably interpret the characteristics of OV in metal oxides, the quantitative study of OV (e.g., content, types, sizes, and configuration) is

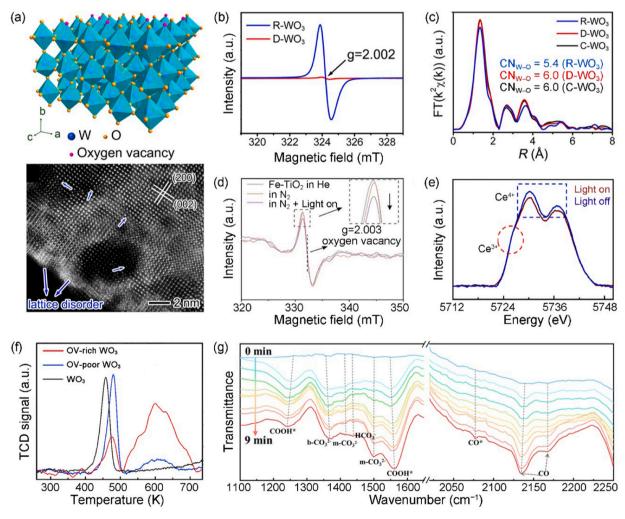


Fig. 7. (a) Schematic illustration of the location of OV in the WO₃ lattice and the corresponding aberration-corrected HAADF-STEM image of defect-rich WO₃ (R-WO₃). (b) Room temperature EPR spectra of R-WO₃ and defect-deficient WO₃ (D-WO₃). (c) Fourier-transform W L₃-edge EXAFS spectra of R-WO₃, D-WO₃, and commercial WO₃ (C-WO₃). Reproduced with permission from Ref. [63]. Copyright 2016, American Chemical Society. (d) In situ EPR spectra of Fe-doped TiO₂ (Fe-TiO₂) before and after light irradiation in the N₂ atmosphere. Reproduced with permission from Ref. [105]. Copyright 2021, John Wiley & Sons, Inc. (e) In situ Ce L₃-edge EAXFS spectra of NiO/CeO₂/rGO before and after light irradiation. Reproduced with permission from Ref. [126]. Copyright 2020, Elsevier B.V. (f) CO₂ TPD spectra of WO₃, OV-poor WO₃, and OV-rich WO₃. Reproduced with permission from Ref. [137]. Copyright 2018, Elsevier B.V. (g) In situ FTIR spectra of photocatalytic CO₂ conversion over oxygen-defective Bi₂O₂CO₃ nanosheets. Reproduced with permission from Ref. [138]. Copyright 2020, John Wiley & Sons, Inc.

essential [127]. PAS is a highly sensitive technique to quantitatively characterize OV-related defects in metal oxides [128]. The lifetime of positrons can shed light on the size, type, and relative concentration of various defects at even the ppm level. Li and co-workers successfully identified two types of OVs in nanosized ceria [129], small neutral Ce3+-OV complex and large-size clusters of vacancy defects. Based on this insight, the authors were able to establish the relationship between the concentration of defect clusters and the reducibility and reactivity of the nanosized ceria. Furthermore, similar defect types (i.e., small neutral Ti³⁺-OV associates, large-size vacancy clusters, and a few voids of vacancy associates) were also observed in hydrogenation-modified TiO2 using PAS [130]. In addition to the defect types and sizes, Zhang et al. [131] recently developed time-resolved optical spectroscopy to "visualize" OV configuration in metal oxides (e.g., WO₃, WO_{2.9}, and WO_{2.72}). Due to the ultrafast characteristic (femtosecond scale) of this technique, it can accurately reflect the information of the excited states, energy bands, and electron transition routes, thus enabling the fingerprint identification of specific OV configurations in metal oxides.

3.2.3. Atomic-scale electron microscopy and scanning tunneling microscopy

In addition to adopting spectroscopic techniques, OV on some model metal oxides (atomic level flatness and clean surface) can be directly imaged by electron microscopy. Through direct atomic-scale TEM observations and density functional theory calculations, Wang and coworkers demonstrated the proton-dissolution-promoted formation, migration, and clustering of cation and anion vacancies during a watervapor-enhanced NiCr alloy oxidation process [132]. Scanning tunneling microscopy (STM) is also commonly used to study OV and its interaction with adsorbed molecules (e.g., H₂O, O₂, and CO₂) [133,134]. It has been successfully employed to image the OV on single crystalline TiO2 under ultra-high vacuum conditions. Schaub et al. observed the oxygen-mediated OV diffusion on the TiO2(110) surface using STM [135]. It was reported that the adsorbed O2 molecules could donate an oxygen atom to OV and then capture an oxygen atom from a neighboring bridging oxygen row, resulting in the observed OV diffusion. The adsorption process of O2 was further studied by a subsequent work using anatase TiO₂(101) as a model surface [136]. It was revealed that the O₂ adsorbed initially on the fivefold-coordinated Ti sites could react with OV to form an $(O_2)_O$ interstitial at the OV site.

3.2.4. CO₂ temperature-programmed desorption and infrared spectroscopy

The highly reactive characteristic of OV makes it an excellent active

site to interact with CO₂ (i.e., CO₂ adsorption and activation). The adsorption of CO2 can be quantitatively studied by CO2 adsorption isotherms and temperature-programmed desorption (TPD). Notably, CO2 TPD can shed light on the adsorption strength of CO2 (i.e., physical adsorption, weak chemical adsorption, and strong chemical adsorption). Liang et al. recently employed CO2 TPD to study the adsorption behaviors of CO₂ on WO₃, OV-poor WO₃, and OV-rich WO₃ [137]. Two peaks at around 460 K and 600 K were observed in TPD spectra and the peak intensity at 600 K was remarkedly enhanced with the introduction of OV, confirming the OV-enabled stronger CO2 adsorption (Fig. 7f). After adsorption, the subsequent activation and conversion of CO2 can produce abundant surface carbon-containing intermediates (e.g., carboxyl group, carbonate, bicarbonate, and CO), which can be monitored by Fourier-transform infrared (FTIR) spectroscopy. For instance, Zu et al. recently applied in situ FTIR to gain insight into the photocatalytic CO2 reduction process over oxygen-defective Bi2O2CO3 nanosheets [138]. Upon light irradiation, two new peaks at 1565 and 1230 cm⁻¹, which can be ascribed to the vibrations of *COOH groups,

Despite the excellent methods described above, there are still some discrepancies to be noted in OV characterization. For example, in the

appeared in the FTIR spectra (Fig. 7g). Moreover, the signal attributed to

adsorbed CO gradually intensified with prolonged irradiation time.

These results collectively suggest a *COOH-mediated CO2-to-CO con-

version pathway on the OV-containing catalyst surface.

EPR spectra, the signal at g = 2.002 is ascribed to single electrontrapped OV sites. Resultingly, the intensity of this signal would be determined by both the concentration and electronic structure of OV. As such, more attention should be taken when analyzing the changes in the signal intensity under light irradiation or during in situ reactions. In the case of XPS, the peak at ~531.0 eV in the O 1s XPS spectra is often ascribed to OV. However, this peak can also be ascribed to surfaceadsorbed oxygen species (e.g., O_2^{2-} and hydroxyl groups). To address this, PL may be an ideal alternative to O 1s XPS for distinguishing and quantifying the OV-induced defect state in some cases (e.g., ZnO-based materials) [139-141]. Despite these shortcomings, it has been well recognized that the EPR and XPS are highly sensitive to the metal valence states, which can indirectly shed some light on the state of OV in metal oxides. Finally, it is highly recommended to combine multiple characterization tools to gain a comprehensive insight into the structure and role of OV.

4. Recent progress of OV engineering in metal oxides for photocatalytic CO₂ reduction

The rapid development of controlled synthesis and advanced characterization techniques for OV-containing metal oxides has laid a foundation for their application in photocatalytic CO₂ reduction. The introduction of OV in metal oxide-based photocatalysts can potentially influence all three basic steps of photocatalytic CO₂ reduction: light absorption, charge separation, and surface CO₂ conversion reaction [142]. The role of OV can differ according to its specific state (e.g., concentration, location, and metal oxide components). For example, a high OV concentration mainly contributes to enhanced light absorption, while a moderate amount of OV on the photocatalyst surface can facilitate charge carrier separation and surface CO₂ conversion. A comprehensive understanding of the role of OV in photocatalytic CO₂ reduction is thus the premise for the rational design of high-performance metal oxide-based photocatalysts.

4.1. Oxygen vacancy-extended light absorption

To meet the thermodynamic requirements for photocatalytic CO_2 reduction with $\mathrm{H}_2\mathrm{O}$, the commonly used metal oxide photocatalysts typically have a wide bandgap (> 3.0 eV). As a result, these metal oxides can only be excited by UV light, significantly limiting their photocatalytic activity toward CO_2 reduction. With the introduction of OV, this limitation can be largely alleviated by extending the light response of metal oxide to the visible light range.

4.1.1. OV-modified band structure

OV can extend the light-absorption range of metal oxide photocatalysts by modifying their band structures (e.g., band positions and defect states) [143]. A typical case is TiO₂, where the addition of OV in commercial white TiO₂ endows it with a variety of colors [77,144], including black, gray, green, blue, and red. The appearance of these colors indicates extended light absorption of metal oxide photocatalysts and is generally accompanied by their enhanced performance toward visible-light-driven CO2 reduction [145-147]. Gao et al. prepared a black titania nanotube array (B-TiO2) by aluminothermic reduction of anodized TiO₂ nanotube arrays [80]. Ascribing to the significantly extended light absorption range, the optimized B-TiO2 catalyst demonstrated an excellent CO2-to-CO conversion yield of 185.4 μ mol g $^{-1}$ h $^{-1}$ under visible light irradiation. In addition, an enhanced CH₄ production rate of 16.2 $\mu mol\ g^{-1}\ h^{-1}$ was observed on hydrogenated blue titania $(H-TiO_{2-x})$ under full solar irradiation. The $H-TiO_{2-x}$ prepared by Li-EDA reduction possessed a crystalline core-amorphous shell structure (TiO2 @TiO2-x) with abundant OVs, thus exhibiting a significantly enhanced light absorption capability [77]. Due to this enhanced light absorption, these OV-containing metal oxides can also be employed to construct hybrid nanostructures for achieving better photocatalytic ${\rm CO_2}$

reduction performance [148,149].

Rationally tuning the OV concentration in metal oxides can enable NIR and even IR light-driven CO₂ reduction reactions. In principle, the bandgap of the photocatalyst should be larger than 1.35 eV to fulfill the CO2-to-CO reduction (-0.12 V vs. NHE) and H2O oxidation to O2 (+1.23 V vs. NHE). Therefore, IR light with wavelengths larger than 920 nm should be incapable of initiating these two reactions simultaneously. To tackle this, Xie and co-workers [137] recently proposed an intermediate band engineering strategy by introducing OV into ultrathin WO₃ atomic layers (Fig. 8a). Theoretical calculations and experimental observations confirmed the formation of an intermediate band in the bandgap of WO3 when the OV concentration reached a critical value, which endows the sample with light responses extending from visible to IR light (up to 10,000 nm) (Fig. 8b). As a result, the OV-rich WO₃ atomic layers showed obvious quantum efficiencies even up to 1064 nm (Fig. 8c). Following this line of thought, Jiang et al. developed a Z-scheme heterojunction by perovskite CsPbBr3 and oxygen-deficient WO₃ for NIR light-driven CO₂ reduction [150]. Under NIR light (780-2500 nm) irradiation, electrons stored in the OV state of WO₃ could be excited to the conduction band of WO₃ and then transferred to the CsPbBr₃ site for CO₂ reduction, achieving CO and CH₄ production rates of 69 μ mol g⁻¹ h⁻¹ and 8 μ mol g⁻¹ h⁻¹, respectively.

Concomitant with the extended light absorption, OV as a donor-like defect state can drive up the Fermi level of metal oxide photocatalysts [151], thus potentially modifying their energy band positions for efficient $\rm CO_2$ reduction. Zhu et al. explored the influence of OV on the photocatalytic $\rm CO_2$ reduction performance of $\rm In_2O_3$ catalysts [86]. It was revealed that the introduction of OV could optimize the conduction band position of $\rm In_2O_3$, allowing photogenerated electrons to be more energetic to participate in $\rm CO_2$ -to-CO conversion. A similar phenomenon was observed in a $\rm Cu^+$ -doped $\rm W_{18}O_{49}$ catalyst [152], which exhibited $\rm \sim 4$ times enhancement in photocatalytic $\rm CO_2$ reduction performance for $\rm CH_4$ production. It was reported that the $\rm Cu^+$ doping could introduce OV and thus change the conduction band edge of $\rm W_{18}O_{49}$ to a more negative position.

4.1.2. OV-induced plasmonic effect

In addition to modifying the energy band structure, OV can also extend the light absorption of metal oxides by inducing a unique plasmonic effect. In some heavily OV-doped metal oxides, the large number of electrons accumulated in the OV-induced defect band can absorb incident light and experience intra-band excitation. Interestingly, such an excitation can generate highly energetic "hot" electrons as well as a significant photothermal effect [66], both of which will contribute to the photocatalytic $\rm CO_2$ reduction process. For example, Zhang and co-workers studied the performance of OV-enriched blue $\rm WO_{3-x}$ porous nanorods for photothermal $\rm CO_2$ conversion [79]. It turns out that the OV enrichment allowed the sample to exhibit enhanced light absorption in the $\rm UV$ -vis-NIR range due to the LSPR effect. As a result, the OV-enriched sample achieved a remarkably enhanced $\rm CH_4$ production rate of $\sim 45.7~\mu mol~g^{-1}~h^{-1}$ under solar light irradiation in the absence of any noble-metal co-catalysts or sacrificial agents.

Recently, Li and co-workers [153] developed an oxygen-deficient Bi₂O_{3-x} catalyst, which exhibited LSPR absorption in the wavelength range of 600–1400 nm (Fig. 8d,e). Due to the enhanced light absorption, the catalyst displayed an apparent quantum yield of 0.113% for CO production at 940 nm (in the presence of H₂ reductant), which was 4-times higher than that observed at 450 nm. Moreover, the catalytic activity was well maintained across three cyclic tests, with each cycle lasting for 8 h, demonstrating the excellent stability of the plasmonic Bi₂O_{3-x} catalyst (Fig. 8f). In another report, Lou and co-workers systematically studied the influence of OV location in Bi₂WO₆ (i.e., W-OV-W and Bi-OV-Bi sites) in inducing the LSPR effect for photocatalytic CO₂ reduction [154]. It was demonstrated that the W-OV-W site could induce an energy state close to the conduction band of Bi₂WO₆, facilitating photoelectron accumulation for a long lifetime. Therefore, a strong LSPR absorption in the 500-1400 nm region was observed for Bi₂WO₆ with W-OV-W sites, which is absent for Bi₂WO₆ with Bi-OV-Bi sites. Benefiting from the LSPR-induced high-energy hot electrons, the Bi₂WO₆ with W-OV-W sites exhibited a CH₄ production rate of 9.95 μmol g⁻¹ h⁻¹, which is 26-fold higher than that of Bi₂WO₆

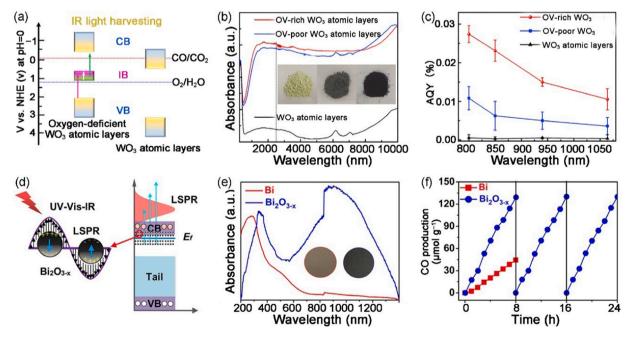


Fig. 8. (a) Schematic illustration for the electronic band structures of WO₃ and oxygen-deficient WO₃ (V_O-rich WO₃) atomic layers. VB, CB, and IB represent the valence band, conduction band, and intermediate band, respectively. (b) Optical absorption spectra of V_O-rich and V_O-poor WO₃ atomic layers. (c) Apparent quantum efficiency measurements for photocatalytic CO₂ reduction over WO₃, V_O-rich, and V_O-poor WO₃ atomic layers. Reproduced with permission from Ref. [137]. Copyright 2018, Elsevier B.V. (d) Schematic illustration for the LSPR excitation process on Bi₂O_{3-x}. (e) UV-vis-IR absorption spectra of Bi and Bi₂O_{3-x} fright). (f) Photocatalytic CO₂-to-CO conversion performance test for Bi and Bi₂O_{3-x} samples. Reproduced with permission from Ref. [153]. Copyright 2020, John Wiley & Sons, Inc.

with Bi–OV–Bi sites (0.37 μ mol g⁻¹ h⁻¹).

The construction of heterojunction is effective in stabilizing the LSPR excitation of metal oxides by maintaining a high free-electron density under light irradiation [155,156], thereby facilitating LSPR-mediated CO₂ reduction [157]. Li and co-worked constructed a heterostructure between TiO2 mesocrystals and plasmonic WO3-x nanowires for co-catalyst-free photocatalytic CO2 reduction [158]. The TiO₂/WO_{3-x} hybrid achieved a much higher CH₄ production rate (16.3 μ mol g⁻¹ h⁻¹) than TiO₂ (3.5 μ mol g⁻¹ h⁻¹) and WO_{3-x} (8.0 μ mol g⁻¹ h⁻¹) under UV-vis light irradiation. Mechanism studies suggested that the continuous photoelectron injection from TiO2 to WO3-x increased the free carrier density of WO_{3-x}, resulting in more stable LSPR absorption and hot electron generation. Similarly, a heterojunction by WO_{3-x} and MoO_{3-x} with abundant OVs was prepared, which displayed significantly enhanced hot electron generation and thus CO2-to-CO conversion performance under UV-vis-NIR light irradiation [159]. In another recent report, Lou and co-workers further demonstrated the plasmonic thermal effect-enhanced CO2 reduction by coupling plasmonic WO_{3-x} with pyroelectric black phosphorous (BP) [160]. Under visible and NIR light irradiation, the local high temperature brought about by the plasmonic thermal effect of WO_{3-x} triggered the pyroelectric effect of BP, inducing the continuous electron injection from pyroelectric BP to WO_{3-x}. As a result, the plasmonic BP/WO_{3-x} hybrid exhibited a CO production rate of 26.1 µmol g⁻¹ h⁻¹, 7- and 17-fold higher than those of WO_{3-x} and BP, respectively.

In short, the introduction of OV can significantly extend the light absorption range of metal oxides by tuning their band structures (i.e., band positions and defect states) and inducing a plasmonic effect. It is worth mentioning that the extended light absorption range does not necessarily lead to enhanced photocatalytic activity. Taking black and gray TiO₂ as an example, black TiO₂ has a broader light absorption range than gray TiO2; however, the latter generally features higher intrinsic photocatalytic activity (in the absence of co-catalysts). In fact, the mismatch between the performance and light absorption capability of the photocatalyst is within expectations because light absorption is only one of the three basic steps (i.e., light absorption, charge carrier separation, and surface reaction) that determine photocatalytic performance. In addition, in OV-containing plasmonic metal oxides, the LSPR absorption is due to the intra-band excitation of electrons in the OVinduced defect band. Therefore, both OV-induced inter-band and intra-band excitations exist in plasmonic metal oxides (e.g., WO_{3-x}) and will potentially affect the catalytic reaction. However, a high electron concentration in the OV-induced defect band is required for maintaining a strong intra-band excitation [154]. To this end, the photocatalytic system should have i) sufficient energy input (e.g., UV light) for electron generation and ii) an efficient charge transfer channel for electron accumulation in the OV-induced defect band. Only after fulfilling these criteria, the OV-induced plasmonic effect can contribute to catalysis by providing highly energetic electrons (hot electron effect) and/or local high temperatures (i.e., photothermal effect).

4.2. Oxygen vacancy-enhanced charge separation

As a donor-like defect, OV can introduce a defect band near the conduction band of metal oxide photocatalysts. The defect band brought about by OV can serve as a trapping site for photogenerated electrons in the conduction band, thereby alleviating the unfavorable recombination of photogenerated charge carriers and facilitating their separation. Taking this into account, numerous researchers have successfully employed OV to maneuver the charge transfer behaviors in pristine metal oxide and metal oxide-based hybrid nanostructures for efficient ${\rm CO}_2$ reduction.

4.2.1. OV in mono-component metal oxides

The role of OV in promoting charge carrier separation has been well-studied by previous works. Wheeler et al. systematically studied the role

of OV in introducing a mid-bandgap state in hydrogen-treated TiO $_2$ (H: TiO $_2$) using ultrafast transient adsorption [161]. Under UV light excitation, the charge carrier relaxation in H:TiO $_2$ was slower than untreated TiO $_2$, indicating the prolonged charge carrier lifetime. In another work, Zhang et al. used EPR to study the light excitation behaviors of OV-rich nanoporous TiO $_{2-x}$ that contains small anatase nanoparticles with high surface area [162]. It was observed that the TiO $_{2-x}$ nanoparticles displayed shallow charge traps associated with Ti/O-vacancies, which remained separated and stable for several minutes after light exposure. The long-lived charge separation state allowed TiO $_{2-x}$ to exhibit a favorable CH $_4$ production rate of 26.12 μ mol g $^{-1}$ h $^{-1}$.

Over time, it has been realized that the function of OV in influencing charge carrier behaviors is largely determined by its locations [163]. A typical case is for bulk and surface OVs, where the latter generally has higher stability [95]. Kong et al. reported a markedly improved separation efficiency of photogenerated electrons and holes with the decreasing relative concentration of bulk OV to surface OV in TiO2 nanocrystalline [163]. Furthermore (Fig. 9a), Yang and co-workers [164] studied the role of bulk and surface OVs in photocatalytic CO₂ reduction and observed an enhanced photocatalytic performance with an increase in the ratio of surface OV (SO) to bulk OV (i.e., single-electron-trapped oxygen vacancy, SETOV). It was revealed that the OV in the surface disorder layer (i.e., surface OV) could promote charge carrier separation, while bulk OV served as recombination centers for photogenerated electrons and holes. However, both bulk and surface OVs can contribute to enhanced visible light absorption of TiO2. Therefore, the TiO₂ sample with the co-existence of surface and bulk OV (TiO₂-SBO) achieved the highest photocatalytic CO₂ reduction performance for CH₄ production under both UV and visible light irradiation (Fig. 9b). It is worth noting that the enhanced light absorption and the inhibited charge carrier separation caused by bulk OV will jointly determine whether there is a final positive or negative effect on the photocatalytic CO2 reduction performance, especially when under UV light irradiation.

The synergy of OV with other surface sites (e.g., hydroxyl groups and metal vacancies) can further promote charge carrier separation and enhance photocatalytic performance [165]. Xiao et al. successfully modified OV and surface hydroxyl groups onto the TiO2 surface (TiO_{2-x}·2xOH), where the OV and the hydroxyl groups with high electron density were able to trap electrons and holes, respectively [166]. The TiO_{2-x}·2xOH sample was synthesized by high-temperature calcination of TiO₂ in the presence of ionic liquid, where the breaking of the Ti-O-Ti bond during high-temperature treatment leads to the formation of OV and surface hydroxyl groups on nearby Ti atoms (TiO₂ + x H₂O \rightarrow TiO_{2-x}·2xOH). Through femtosecond transient absorption measurements, Ozin and co-workers also observed a longer excited-state lifetime in the In₂O₃ catalyst co-modified with OV and hydroxyl groups (In₂O_{3-x}(OH)_v) [167]. This observation was a result of improved charge separation due to the defect states induced by OV (shallow donor state, electron trapping) and hydroxyl groups (shallow acceptor state, hole trapping) in the bandgap of $In_2O_{3-x}(OH)$. As a result, the $In_2O_{3-x}(OH)_v$ with OV and hydroxyl groups on its surface functioned as an effective gas-solid-phase photocatalyst for CO2-to-CO reduction in the presence of H₂.

In addition to serving as an electron trapping site, OV can also indirectly contribute to the charge carrier separation process by assisting in the formation/stabilization of the electric field within the metal oxide photocatalyst. Recently, Ma and co-workers [168] reported the synergy of OV and ferroelectric polarization in accelerating charge carrier separation and promoting $\rm CO_2$ photoreduction on oxygen-deficient polarized $\rm Bi_3NbTiO_9$ nanosheets (BNT-OVP) (Fig. 9c). Interestingly, the authors observed a pinning effect of OV on the ferroelectric domains of BNT-OVP. In other words, the strong electric field induced by the superb ferroelectric polarization could be well maintained by OV to facilitate the charge carrier separation within the BNT-OVP nanosheets. Compared to OV-deficient polarized $\rm Bi_3NbTiO_9$ nanosheets (BNT-P), the

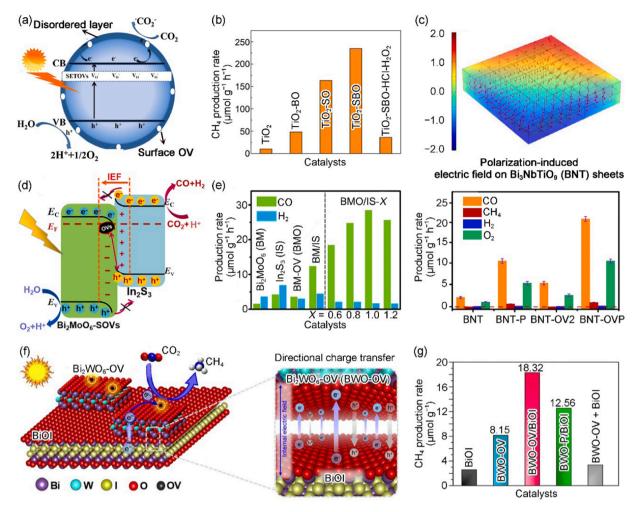


Fig. 9. (a) Schematic illustration for the photocatalytic CO₂ reduction process over the TiO₂ sample (TiO₂-SBO) containing both surface and bulk OVs. (b) Photocatalytic CH₄ production rates over different TiO₂ samples under UV light irradiation. Reproduced with permission from Ref. [164]. Copyright 2017, Elsevier B.V. (c) COMSOL simulation of polarization-induced electric field on polarized Bi₃NbTiO₉ (BNT-P) nanosheets. And the photocatalytic production rates of CO, CH₄, and H₂ over BNT, BNT-P, oxygen-deficient Bi₃NbTiO₉ (BNT-OV), and polarized BNT-OV (BNT-OVP). Reproduced with permission from Ref. [168]. Copyright 2021, Springer Nature. (d) Schematic illustration for the charge transfer mechanism over the Z-scheme oxygen-deficient Bi₂MoO₆/In₂S₃ heterojunctions. IEF: internal electric field. (e) Photocatalytic CO and H₂ production rates over different samples. Reproduced with permission from Ref. [180]. Copyright 2021, Elsevier B.V. (f) Schematic illustration for the OV-directed charge transfer from BiOI to OV-rich Bi₂WO₆ (BWO-OV). (g) Photocatalytic CH₄ production rate over BiOI, BWO-OV, BWO-OV/BiOI heterojunction, OV-deficient heterojunction (BWO-P/BiOI), and physical mixture of BWO-OV and BiOI. Reproduced with permission from Ref. [183]. Copyright 2021, Elsevier B.V.

BNT-OVP demonstrated nearly a 2-fold enhancement in CO (20.9 $\mu mol~g^{-1}~h^{-1})$ and CH₄ ($\sim\!10.0~\mu mol~g^{-1}~h^{-1})$ production under AM 1.5 G light irradiation.

4.2.2. OV in metal oxide-based hybrid nanostructures

The introduction of OV brings about more complex but fascinating influences on the charge transfer/separation behaviors in metal oxide-based hybrid nanostructures. Theoretically, an increase in the density of electron donor-like defect states (e.g., OV) can lead to a decrease in the work function of metal oxides [151]. With this in mind, Zheng et al. [169] successfully utilized OV to tune the energy match level between $\rm TiO_2$ and polyoxometallate (POM), achieving multi-electron transfer from the OV-rich $\rm TiO_2$ to POM ($\rm H_3PW_{12}O_{40}$, HPW). With this supplement of electrons, the 5 wt% HPW-modified defective $\rm TiO_2$ exhibited the highest selectivity (still < 4%) for CH $_4$ production. Furthermore, Chong et al. introduced a MgAl layered double oxide (MgAl-LDO) into the interlayer between Pt nanoparticles and $\rm TiO_2$ substrate to mediate electron transfer within the hybrid [170]. The OV arising from MgAl-LDO deposition could suppress the recombination of photogenerated charge carriers, enabling multi-electron transfer from $\rm TiO_2$ to

Pt sites. As a result, the MgAl-LDO introduction enhanced the CO and CH_4 yield by 2 and 11 times, respectively, with a greatly enhanced selectivity toward CH_4 production.

The role of OV in mediating the electron transfer pathway has also been observed in OV-containing Z-scheme heterojunctions [171,172]. Z-scheme heterojunction can facilitate the separation of photogenerated electrons and holes while retaining their high redox capabilities [173-175]. Therefore, the introduction of OV to assist or induce Z-scheme charge transfer is highly promising for constructing high-performance photocatalysts for CO_2 reduction [176–179]. Yu et al. recently developed a hierarchical Bi_2MoO_6/In_2S_3 heterostructure with abundant OVs for efficient CO2 photoreduction [180]. Theoretical calculations and in situ XPS observations confirmed a Z-scheme electron transfer pathway (Fig. 9d), where energetic electrons and holes could remain in the conduction band of In₂S₃ and valence band of Bi₂MoO₆, respectively, to participate in photocatalytic CO₂ reduction with H₂O. Based on experimental results, the authors further proposed that the OV on Bi₂MoO₆ can accelerate the Z-scheme electron transfer from Bi₂MoO₆ to In₂S₃. As such, the optimized nanohybrid demonstrated a much higher yield of CO (28.5 µmol g⁻¹ h⁻¹) than that of defect-free

Bi $_2$ MoO $_6$ /In $_2$ S $_3$ heterojunction (\sim 12.0 μ mol g $^{-1}$ h $^{-1}$) (Fig. 9e). In another study of H $_2$ O oxidation using the Z-scheme heterojunction composed of BiOI and OV-enriched CeO $_2$, Sultana et al. demonstrated that the OV could mediate the Z-scheme electron transfer from the conduction band of CeO $_2$ to the valence band of BiOI [181].

In addition to mediating electron transfer within the hybrid nanostructures, OV can also direct the electron transfer pathway for promoted charge separation [182-185]. Using Na-EDA as a selective reductant, Hwang et al. developed a heterophase junction between disordered anatase and ordered rutile (A_d/R_o) TiO₂ [186]. The OV in the disordered anatase not only extends the light absorption range of the hybrid but also serves as an electron trapping site to drive the electron transfer from R_o to A_d . As a result, the A_d/R_o TiO_2 displayed the highest CH₄ production rates of 3.983 µmol g⁻¹ h⁻¹ under visible light irradiation, which was even better than that of metal co-catalyst loaded commercial P25. Kong et al. developed a 2D/2D heterojunction between BiOI and OV-rich Bi₂WO₆, in which the abundant OVs in Bi₂WO₆ induced a directional electron transfer from BiOI to Bi₂WO₆ (Fig. 9f) [183]. As a result, the OV-rich heterojunction exhibited an enhanced CH₄ production rate (18.32 µmol g⁻¹) compared to the OV-deficient $Bi_2WO_6/BiOI$ heterojunction (12.56 µmol g⁻¹) (Fig. 9g). Similarly [187,188], Miao et al. developed a co-catalyst BiO_{2-x} coupled g-C₃N₄ photocatalyst [189], which showed a CO production rate (42.29 µmol g_{cat}^{-1}) 3.8 times that of pure g-C₃N₄. Sun et al. also developed a highly efficient photocatalyst consisting of reduced graphene oxide (rGO) and BiO quantum dots (QDs) [190]. The apparent quantum efficiency of rGO-BiO QDs reached 4.2% at 300 nm, which was ~15 times that of rGO-TiO₂ QDs (0.28%) under the same conditions. Mechanism studies suggested the role of OV in inducing a unique circular electron transfer pathway, i.e., from the conduction band of BiO to rGO and then to being trapped by the OV site on BiO. In other words, electrons were stored in rGO before participating in photocatalytic CO₂ reduction at the OV site.

In short, the electron-trapping effect of OV makes it a useful tool for regulating charge separation behaviors in metal oxide-based photocatalysts (Fig. 10). Tremendous efforts have been devoted to understanding the influence of the state of OV (e.g., locations) on the charge separation process. Despite the deeper understanding obtained, there are still some areas of contention that need to be clarified in future works. For example, in the OV-containing heterojunctions, it is commonly believed that OV at the interface can mediate the electron transfer between the two components. However, the interfacial OV can also be intuitively regarded as the bulk OV, which is supposed to be unfavorable for the charge separation process. For example, Zhang et al. reported that the OV at the interface between TiO₂ and perovskite can slow down charge transfer and significantly accelerate charge recombination at the interface [191]. A rational photocatalyst design with controllable ratios of interface OV to surface OV is hence necessary to

unravel the role of OV in metal oxide-based heterojunction photocatalysts.

4.3. OV-promoted surface CO₂ conversion reaction

The introduction of OV via the removal of oxygen atoms from the metal oxide matrix is generally accompanied by the formation of coordinatively unsaturated metal centers (M_{CUS}). In principle, these M_{CUS} –OV pairs can serve as ideal sites for molecular (e.g., CO_2) chemisorption. Furthermore, the ability of the neutral M_{CUS} –OV pair to accumulate electrons under light irradiation allows them to be an efficient channel for electron transfer between the metal oxide and adsorbed CO_2 molecules. In addition, the OV can potentially modify the electronic structure of other active sites on the metal oxide surface, thus indirectly contributing to the surface CO_2 conversion process. Hence, promoted surface CO_2 conversion reaction is commonly observed on OV-containing metal oxide photocatalysts.

4.3.1. OV as the active site for CO₂ adsorption and activation

The adsorption of CO₂ on the catalyst surface is indispensable for its activation and conversion. Generally, CO2 can be adsorbed with a configuration of carbon coordination, oxygen coordination, or mixed coordination (i.e., using both carbon and oxygen atoms) [6]. The presence of OV on metal oxide-based photocatalysts has significant influences on the adsorption behaviors of CO₂ [58,192,193]. Sorescu et al. employed dispersion-corrected density functional theory to study the adsorption of CO₂ on the anatase TiO₂(101) surface with bridging OVs [194]. It was reported that the most stable configuration of CO₂ adsorption was a bent CO2 structure bound to two surface Ti atoms adjacent to the bridging OV. Moreover, subsurface OVs were found to further enhance the binding of CO2 molecules to the surface. The OV-enhanced CO2 adsorption not only accelerates the photocatalytic reaction [97,195,196] but also allows the reaction to be carried out at low CO2 concentrations. Chen et al. recently employed OV-rich NiO nanoplatelets (r-NiO) to reduce diluted CO2 (10 vol%) under visible light irradiation (>400 nm) [89]. The r-NiO nanoplatelets were synthesized by calcination under Ar and then integrated with a Ru-based photosensitizer, enabling visible-light-driven CO2 reduction (in the presence of TEOA as a hole sacrificial agent). Benefiting from the strong adsorption of CO2 on OV sites, the r-NiO exhibited nearly a 2-fold enhancement in the production rate (6.28 mmol g⁻¹ h⁻¹) of CO from diluted CO₂ reduction compared to that of normal NiO (3.94 mmol g⁻¹ h^{-1}). The role of OV in promoting diluted CO₂ reduction is vital for its practical application.

When CO₂ is adsorbed at the OV site, it can accept electrons from adjacent metal atoms [197] and undergo activation and dissociation for its subsequent transformation [198–200]. Lee et al. [201] employed

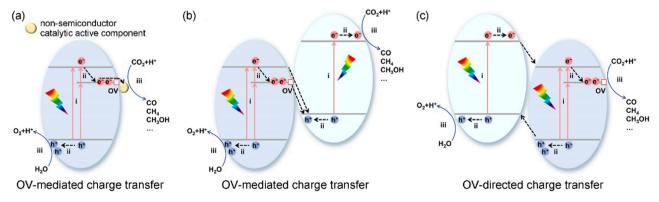


Fig. 10. Schematic illustration for OV-mediated charge transfer (a) between metal oxides and non-semiconductor catalytic active components (e.g., metal nanoparticles) and (b) within semiconductor heterojunction. (c) Schematic illustration of OV-directed charge transfer. The energy band bending at the contact interface and the recombination process of photogenerated electrons and holes are not shown in the figure.

STM to investigate the dissociation of CO_2 adsorbed at the OV site on the $TiO_2(110)$ surface at the single-molecular level (Fig. 11a,b). It was demonstrated that the electron injection from the STM tip into the adsorbed CO_2 induced the dissociation of CO_2 , which was accompanied by the healing of the OV site. The intermediate negative ion ($\bullet CO_2$) was considered to be essential for CO_2 activation and dissociation. In a subsequent theoretical study [202], Chu et al. reported the formation of intermediate $\bullet CO_2$ at the OV site on rutile $TiO_2(110)$ by single electron transfer from the conduction band of TiO_2 to the adsorbed CO_2 molecule. It was shown that if the intermediate $\bullet CO_2$ possessed a lifetime longer than 12 fs, it can excite the bending and asymmetric stretching vibrations with the assistance of OV (Fig. 11c). During this process, the CO_2 LUMO could be sufficiently stabilized below the conduction band minimum of TiO_2 , allowing it to trap photoexcited electrons and dissociate

The activation and dissociation processes of CO_2 on OV have also been experimentally evidenced in recent reports [203,204]. Wang et al. reported the dissociation of CO_2 into carbon and oxygen over an OV-rich amorphous zinc germanate (α -Zn-Ge-O) photocatalyst under xenon lamp irradiation [205]. EPR and isotope labeling studies confirmed the role of OV and photogenerated electrons in promoting CO_2 dissociation; meanwhile, photogenerated holes can reoxidize surface oxygen atoms

into O₂ to regenerate OV on the photocatalyst surface. Similarly, Liu et al. carried out photocatalytic CO2 reduction studies on three TiO2 nanocrystal polymorphs (anatase, rutile, and brookite) that were engineered with defect-free and OV-rich surfaces [206]. The helium treatment of the TiO2 nanocrystals at moderate temperatures created OV and Ti³⁺ on anatase and brookite but not on rutile. In situ diffuse-reflectance infrared Fourier transform spectroscopy (DRIFTS) analyses showed that the •CO₂ intermediate could only be observed on OV-containing samples (i.e., anatase and brookite TiO2), in agreement with the OV-associated single-electron transfer process. Following this line of thought, Xie and co-workers successfully achieved the photofixation of CO2 with methanol to long-chain chemicals (i.e., dimethyl carbonate, DMC) on OV-rich Bi₂O₃ nanosheets (Fig. 11d) [207]. Both in situ DRIFTS and XPS revealed that the OV could provide abundant localized electrons for the enhanced generation of •CO₂, which was considered the rate-determining step for CO₂ photofixation. As a result, the CO₂ conversion over OV-rich Bi₂O₃ reached ~18%, 9 times higher than that of OV-poor Bi₂O₃ (Fig. 11e).

In addition to the promoted CO_2 adsorption and activation, OV can also alter the CO_2 conversion pathway. Liu et al. investigated the CO_2 reduction pathway on defective anatase $TiO_2(101)$ surface with density functional theory calculation [208]. It was reported that the OV not only

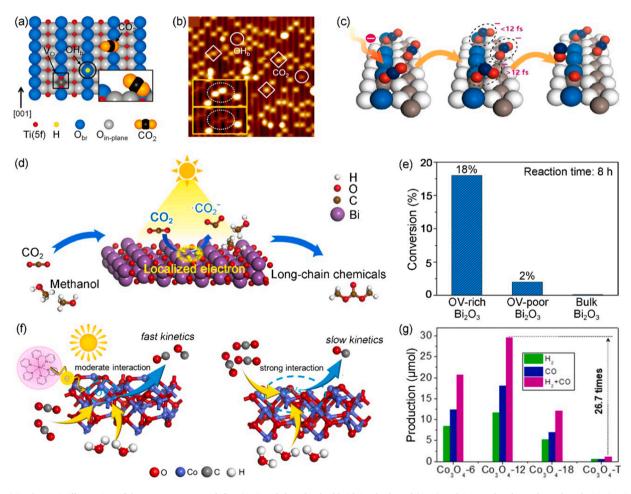


Fig. 11. (a) Schematic illustration of the oxygen vacancy defect (OV) and the adsorbed bridging hydroxyl (OH_b) and CO₂ molecule on the reduced TiO₂(110) surface. (b) STM image of the TiO₂(110) surface after adsorption of CO₂ at 55 K, showing the occupation of all the OV sites by CO₂. CO₂ and OH_b features are highlighted with diamonds and circles, respectively. The insets show the OV site with and without adsorbed CO₂. Reproduced with permission from Ref. [201]. Copyright 2011, American Chemical Society. (c) Schematic illustration for the photoexcited electrons-driven CO₂ activation process on the OV site of rutile TiO₂(110) surface. Reproduced with permission from Ref. [202]. Copyright 2020, American Chemical Society. (d) Schematic illustration for the photocatalytic reduction of CO₂ to long-chain chemicals over Bi₂O₃ nanosheets with rich oxygen vacancies (OV-rich Bi₂O₃). (e) Photocatalytic CO₂ reduction performance of OV-rich Bi₂O₃, OV-poor Bi₂O₃, and bulk Bi₂O₃ at 373 K under xenon lamp irradiation. Reproduced with permission from Ref. [207]. Copyright 2019, Springer Nature. (f) Schematic illustration for photocatalytic CO₂ reduction performance of various samples under visible light irradiation ($\lambda \ge 420$ nm). Reproduced with permission from Ref. [211]. Copyright 2021, Elsevier B.V.

promoted CO2 adsorption and activation but also stabilized CO and other intermediates for the subsequent hydrogenation and production of CH₄. Furthermore, Ji et al. demonstrated that the hydrogenation of surface intermediates could occur at both the surface Ti and the OV sites, while deoxygenation processes largely take place at the OV site [209]. During the reduction of CO₂ over defective anatase TiO₂(101), the CO₂ first occupied the OV site to undergo deoxygenation and generate CO. It subsequently experiences fast hydrogenation to hydrocarbon products. Following this work, the authors further studied the role of H₂O in the photocatalytic CO2 reduction reaction over the defective anatase TiO₂(101) [210]. Interestingly, it was found that H₂O directly donated protons to CO2 and other reaction intermediates adsorbed at the OV site while the remaining OH⁻ recombined with a surface-adsorbed proton to generate a new H₂O molecule. This pathway is distinctly different from the traditional understanding that H₂O first experiences photooxidation to produce protons, which then participate in CO₂ hydrogenation and reduction. However, the possibility of H2O participating in CO2 reduction by combining a proton to form a hydrated proton (e.g., H₃O⁺) should also be duly considered.

Despite the significance of OV in promoting $\rm CO_2$ adsorption, activation, and conversion, rationally tuning the amount of OV is still essential. Zhang et al. [211] reported that the $\rm Co_3O_4$ catalyst decorated with moderate OV amounts exhibited the highest performance for CO production in the presence of TEOA as a sacrificial agent (Fig. 11f,g). It was reported that excessive or insufficient OVs on the $\rm Co_3O_4$ surface would lead to sluggish surface reaction kinetics. Excessive amounts of OVs lead to overly strong interactions between catalysts and substrate molecules, thus blocking the catalytic cycle. On the other hand, the

substrate CO_2 cannot be efficiently activated if the interactions are too weak (i.e., when there are insufficient OVs). In addition to CO_2 adsorption, an excessive amount of OVs can also impede the activation process of CO_2 adsorbed at the OV site. Zhao et al. reported that large amounts of OVs on the anatase TiO_2 nanorod surface led to severe charge carrier recombination, significantly limiting the single-electron transfer associated CO_2 activation process [212]. It was also demonstrated that atomic layer deposition (ALD) could be employed as an effective tool to passivate the excessive surface OV states.

4.3.2. OV-based active site design

During photocatalytic CO_2 reduction, the coordinatively unsaturated metal centers near OV also contribute to the stabilization of CO_2 reduction intermediates, constituting the M_{CUS} –OV pair sites. Yan and co-workers reported the Nb–OV pair site-assisted CO_2 conversion on mesoporous black Nb_2O_{5-x} catalyst [213]. It was shown that the Nb–OV pair site not only assisted the adsorption and dissociation of CO_2 but also facilitated the stabilization of *CO intermediates and their transformation into *CHO (Fig. 12a). Therefore, as compared to defect-free white Nb_2O_5 (0%), the black Nb_2O_{5-x} photocatalyst achieved a higher CH_4 selectivity (64.8%) from CO_2 reduction (Fig. 12b).

Metal modification can also efficiently influence the photocatalytic CO_2 reduction process by modifying the M_{CUS} –OV pair site [65,214]. The introduced metal centers can stabilize the adjacent OV sites [215–217] and, together with the OV, promote CO_2 adsorption and activation [218–220]. For instance, Chen et al. reported the enhanced CO_2 to CH_4 conversion over OV-rich Zn-doped ultrathin COO layer [221]. The Zn–OV pair site promoted the formation of *CO

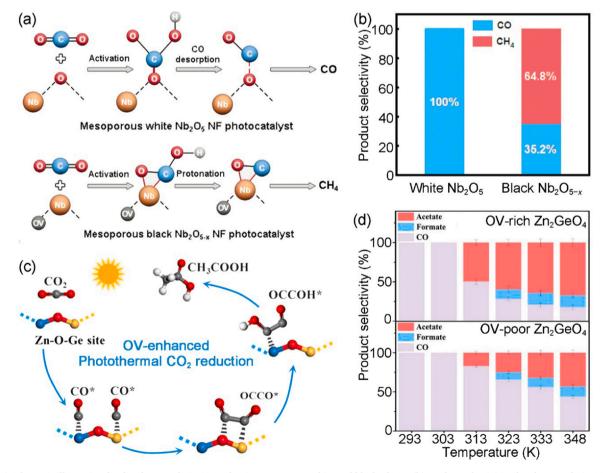


Fig. 12. (a) Schematic illustration for the photocatalytic CO_2 reduction routes over white and black Nb_2O_5 . (b) Product selectivities for photocatalytic CO_2 reduction over white and black Nb_2O_5 . Reproduced with permission from Ref. [213]. Copyright 2022, John Wiley & Sons, Inc. (c) Schematic illustration for the OV-enhanced photocatalytic CO_2 reduction to acetic acid process over Zn_2GeO_4 . (d) Product selectivities for photocatalytic CO_2 reduction over OV-rich and OV-poor Zn_2GeO_4 nanobelts in simulated air at different temperatures. Reproduced with permission from Ref. [239]. Copyright 2021, American Chemical Society.

intermediates while the π -back-bonding between Zn dopant and *CO stimulated the CH₄ evolution. As a result, the OV-rich Zn-CoO layer exhibited enhanced CO₂ reduction efficiency (26.8 μ mol g⁻¹ h⁻¹) and CH₄ selectivity (63.8%) compared to 7.2 μ mol g⁻¹ h⁻¹ and 24.6% for the undoped OV-CoO layer. However, the photocatalyst efficiencies for OV-rich Zn-CoO and undoped OV-CoO reduced to 83.4% and 52.8% during the second catalytic cycle, respectively. This suggests that the OV on the catalyst surface was consumed during the photocatalytic reaction.

The role of metal dopants in stabilizing CO_2 reduction intermediates has also been observed in Cu^+ -modified TiO_2 [222–224]. Liu et al. reported that the Cu^+ dopant could assist in the destabilization and dissociation of the ${}^{\bullet}CO_2^-$ intermediate formed at the adjacent Ti^{3+} –OV site due to the strong interaction between Cu^+ and *CO intermediates [225]. Furthermore, an over 30 times increase in CH_4 production rate (8.68 µmol g^{-1} h^{-1}) was observed over Cu_2O cluster-loaded TiO_2 nanosheets compared to the parent TiO_2 [226]. Extending from this, Wang et al. recently developed an Au/TiO_{2-x} hybrid, where the slightly positively-charged Au formed under plasmonic excitation can stabilize the in situ generated *CO from the OV site [227]. Notably, it was demonstrated that the *CO stabilization allowed for efficient C–C coupling over the Au/TiO_{2-x} hybrid, achieving a 20% selectivity toward C_2H_6 production (0.66 µmol g^{-1} h^{-1}) from photocatalytic CO_2 reduction.

Introducing additional active sites for enhanced H2O activation and proton production is another promising pathway for promoting photocatalytic CO₂ reduction [228-231]. In this regard, Pt has commonly been considered the ideal choice for H₂O activation [232-234]. Yu and co-workers developed a Pt nanoparticles-loaded OV-rich Ga₂O₃ for enhanced photocatalytic CO₂-to-CO conversion [235]. It was reported that the hydrogen atoms formed from H2O reduction at the Pt site could be transferred to the adsorbed CO₂ molecules at the OV site on Ga₂O₃. As a result, the OV-rich Pt/Ga_2O_3 exhibited a substantially enhanced CO production rate (21.0 μmol h⁻¹) compared to that of OV-poor Pt/Ga₂O₃ $(3.9 \, \mu \text{mol h}^{-1})$ under xenon lamp irradiation. In addition, it was recently demonstrated that thermal coupling could facilitate the electron transfer from OV to Pt sites, leading to promoted H2O activation and hydrogen-assisted CO2 reduction. During the thermally coupled photocatalytic CO2 reduction, an enhanced hydrogen splitting was observed over Pt/TiO_{2-x} with smaller Pt nanoparticles (2.33 nm vs. 4.06 nm), which is beneficial for the subsequent CO₂ activation by hydrogenation [236]. Unsurprisingly, a 43-fold increase in CO₂ conversion rates was observed when the system temperature was increased from 298 to 393 K. With that in mind, Cai et al. recently developed an OV-rich Au/TiO2 photocatalyst, where the plasmonic effect of Au can locally provide high temperatures to induce the thermally coupled CO2 reduction over the OV site [237]. The synergistic effect between OV and other active sites on the photocatalyst can potentially lead to higher performance for CO2 reduction. However, the rational integration of the multiple active sites to maximize their synergistic effect remains challenging.

4.3.3. OV-assisted active site design

Apart from acting directly as the active site for CO_2 reduction, OV can also indirectly contribute to CO_2 reduction by affecting the electronic structures of actual catalytic active sites [238]. Recently, Xie and co-workers reported the OV-enhanced photothermal CO_2 reduction for acetate production over OV-rich Zn_2GeO_4 nanobelts [239]. DFT calculations demonstrated that the asymmetric Zn-O-Ge site endowed the two neighboring *CO molecules with different charge densities, which was beneficial for their coupling and hence the production of *OCCO intermediates (Fig. 12c). Moreover, the OV in the material further enhanced the asymmetry of the Zn-O-Ge site, thus leading to improved production of *OCCO and its hydrogenation to *OCCOH intermediates. The OV-rich Zn_2GeO_4 nanobelts displayed an acetate production yield of 12.7 μ mol g^{-1} h^{-1} with a selectivity of 66.9% in a simulated air

atmosphere at 348 K, significantly outperforming that of OV-poor $\rm Zn_2GeO_4$ (production yield: 1.4 $\mu mol~g^{-1}~h^{-1}$, selectivity: 43.1%) (Fig. 12d).

Beyond that, the unique geometry and electronic structures of OV enable it to be an excellent site for the deposition of catalytically active components [240,241]. For instance, the redox reaction between the electrons trapped at the OV sites and the metal salt precursors can result in the deposition of metal nanoparticles on the defective metal oxide support [242]. Following the OV-assisted deposition, the intimate contact between OV-containing metal oxide and the deposited catalytic active components allows the formation of a high-quality interface. Therefore, a high charge transfer efficiency and synergistic effect between the two components can be expected [243], resulting in an augmented photocatalytic CO₂ reduction performance. In this regard, Cai et al. successfully employed OV to anchor single atomic Pt sites on the defective TiO₂ surface, where the synergy of OV and the anchored Pt contributed to enhanced CO₂ reduction performance [244]. Experimental and theoretical studies suggested that the strong electronic interaction between TiO₂ and the OV-anchored Pt atoms can impede the aggregation of Pt. At the same time, Pt modification favored the generation of more surface and subsurface OVs in the TiO2 substrate.

In addition to metal deposition, Li et al. recently developed an OVassisted strategy to construct FeOOH/CdS heterostructure as an efficient photocatalyst for CO2 reduction [245]. It was found that the CdS nanocrystals with a size of ~9.3 nm were preferentially deposited at the OV sites on FeOOH nanospindles, forming a hierarchical type-II heterostructure. Due to the synergistic effect of the excellent interfacial charge transfer and the OV-assisted surface reaction, the optimized OV-rich FeOOH/CdS hybrid demonstrated a remarkable CO production rate of 12.55 $\mu mol~g^{-1}~h^{-1}$ under visible light irradiation, about 5.0 times higher than that of OV-rich FeOOH (2.52 μ mol g⁻¹ h⁻¹). This work highlighted the pivotal role of the OV-assisted synthetic strategy in the formation of high-quality interfaces within the metal oxide-based hybrid nanostructure. However, this OV-assisted synthesis strategy is largely applied to supported metal catalysts. Extending this strategy to the synthesis of more types of metal oxide-based hybrid nanocatalysts will be promising for the discovery of high-performance photocatalysts.

5. Conclusions and future perspectives

OV represents a versatile tool for improving the performance of metal oxide-based photocatalysts for CO2 reduction. In this review, we have summarized the role of OV in metal oxide-based photocatalytic CO₂ reduction (Table 2), together with its controllable synthesis and advanced characterization. Generally, the controllable OV synthesis can be approached from two perspectives, applying a reductive environmental condition and tuning the reducibility of the metal oxide (i.e., OV formation energy). The introduced OV can significantly affect the photocatalytic CO₂ reduction performance of metal oxides by extending the light-absorption range, enhancing the separation of photogenerated charge carriers, and promoting the surface CO2 conversion reaction. Moreover, the detailed working mechanism of OV in affecting these steps is highly dependent on its states, such as concentration, location, and the component of substrate metal oxides. Based on this understanding, one can rationally design OV to achieve specific functions for metal oxide-based photocatalysts. Particularly, the three basic steps in photocatalytic CO₂ reduction can be simultaneously optimized by OV engineering, with the aim of obtaining high-performance metal oxidebased photocatalysts for CO2 reduction.

Although remarkable progress has been made in utilizing OV-containing metal oxide photocatalysts for CO_2 reduction, there are still many considerable challenges that have yet to be addressed. Firstly, the performance of the current OV-containing metal oxide photocatalysts is far from satisfactory for its practical application. Under ambient conditions (room temperature and 1 atm pressure), the current CO_2 conversion is commonly at the μ mol g^{-1} level, which should at least be

Table 2
Summary of the roles of OV in promoting metal oxide-based photocatalytic CO₂ reduction.

^a Role of OV			Catalyst	Product yields	^e Performance	Reaction conditions	Ref.
Light absorption ^b (B, P)	Charge separation ^c (M, H)	Surface reaction ^d (A,D)		$(\mu \text{mol } g^{-1} h^{-1})$	enhancement		
В	_	_	WO ₃ layer	CO, 2.8	CO, 1.8 ×	H ₂ O vapor, 40 W IR lamp (> 800 nm)	[137]
В	-	_	Bi_2WO_6	CH ₄ , 4.3	CH_4 , 3.0 \times	H ₂ O vapor, 300 W Xe lamp (> 400 nm)	[75]
P	_	-	WO_{3-x}	CH ₄ , 45.7	CH ₄ , 45.7 ×	$\rm H_2O$ vapor, 300 W Xe lamp (AM 1.5 G filter)	[79]
P	-	_	Bi_2WO_6	CH ₄ , 9.95	CH_4 , 26.9 \times	H ₂ O vapor, 300 W Xe lamp	[154]
P	_	-	MoS ₂ /MoO _{3-x}	CO, 32.4 CH ₄ , 2.0	CO, 12.9 ×	H ₂ O vapor, 300 W Xe lamp	[157]
P	-	-	WO ₃ /black phosphorus	CO, 26.1	CO, 17.0 ×	$\rm H_2O$ vapor, g 86 °C, Xe lamp (> 400 nm)	[160]
В	M	-	TiO ₂	CH_4 , 235.24 ppm $g^{-1} h^{-1}$	CH ₄ , 22.8 \times	0.1 M KHCO ₃ , 250 W Hg lamp	[164]
-	M	-	Bi ₃ NbTiO ₉ nanosheets	CO, 20.9	CO, 2.0 \times	$NaHCO_3$, 4 M H_2SO_4 , Xe lamp (AM 1.5 G)	[168]
В	M	Α	Sr ₂ Bi ₂ Nb ₂ TiO ₁₂ sheets	CO, 17.11	CO, 5.5 ×	NaHCO ₃ , 4 M H ₂ SO ₄ , 20 °C, Xe lamp	[142]
	Н	-	$Bi_2Mo_2O_{6-x}$ / In_2S_3	CO, 28.54	CO, 2.5 \times	H_2O , 300 W Xe lamp (> 420 nm)	[180]
_	Н	_	Bi ₂ WO _{6-x} /BiOI	CH ₄ , 2.3	CH_4 , 1.5 \times	H ₂ O, 500 W Xe lamp (> 400 nm)	[183]
-	Н	-	rGO/BiO	CO, 4.5 CH ₄ , 21.75	-	$\rm H_2O$, 4 °C, 300 W Xe lamp	[190]
-	-	Α	Co ₃ O ₄	CO, 18.0 µmol h ⁻¹	CO, 29.5 \times	$[Ru(bpy)_3]Cl_2$, H_2O , CH_3CN , ^h TEOA, Xe lamp (≥ 420 nm)	[211]
-	-	Α	r-NiO	CO, 3.14 μmol h ⁻¹	CO, 1.6 \times	[Ru(bpy) ₃]Cl ₂ , H ₂ O, 10% CO ₂ , CH ₃ CN, TEOA, 5 W LED (400–1000 nm)	[89]
-	-	Α	Bi ₂ O ₃ nanosheet	f DMC, 18% conversion	9.0 ×	$\mathrm{CH_{3}OH,100~^{\circ}C,0.2Mpa,300~W}$ Xe lamp	[207]
-	-	D	Black Nb ₂ O ₅	CO, 10.6 CH ₄ , 19.5	S _{CH4} , 0–64.8%	H_2O , 300 W Xe lamp (> 420 nm)	[213]
-	-	D	Zn-doped CoO	CO, 9.7 CH ₄ , 17.1	CO, 1.8 \times CH ₄ , 10 \times	H_2O , 80 Pa CO_2 , Xe lamp (AM 1.5 G filter)	[221]
_	_	D	Cu-doped CeO _{2-x}	CO, 1.65	CO, 26.0 ×	H ₂ O vapor, 300 W Xe lamp	[215]
-	-	D	Au/TiO _{2-x}	CH ₄ , 2.66 C ₂ H ₆ , 0.66	-	H ₂ O vapor, 50 W LED (530 nm)	[227]
-	Н	D	Pt/Ga ₂ O _{3-x}	CO, 105.0 CH ₄ , 4.8	CO, 5.4 ×	$\rm H_2O,300~W~Xe~lamp$	[235]
-	_	D	Au/TiO _{2-x}	CO, 3.4 CH ₄ , 5.0	CO, 2.7 \times CH ₄ , 1.5 \times	$\rm H_2O,^g181^{\circ}C,300$ W Xe lamp	[237]
-	-	D	Zn ₂ GeO ₄ nanobelt	Acetate, 12.7 CO, 13.5 HCOOH, 11.7	Acetate, 9.1 ×	$\rm H_2O,0.03\%CO_2$ in Ar, 75 °C, 300 W Xe lamp (AM 1.5 G filter)	[239]

^a OV can influence all three basic steps in photocatalytic CO₂ reduction. Herein, only the most significant influences are listed for a better comparison.

improved by 3 orders of magnitude to fit the continuous process in industrial manufacturing. The limited catalytic activities also reduce the possibility of carrying out the photocatalytic CO_2 reduction in low concentrations of CO_2 (i.e., atmospheric air). Moreover, the current products from photocatalytic CO_2 reduction are mainly C_1 compounds (i.e., CO, CH_4 , CH_3OH), which are less valuable than multi-carbon compounds (C_{2+} , e.g., C_2H_6 , C_2H_4 , C_2H_5OH , CH_3CHO , and CH_3COOH). Therefore, rational active site design on OV-containing metal oxides for enhanced CO_2 adsorption and efficient C–C coupling is highly demanded. In addition, coupling photocatalytic CO_2 reduction with thermal and/or electric can be a potential choice for further performance optimization.

Secondly, the true performance of the photocatalyst needs to be carefully evaluated. During photocatalytic CO_2 reduction, any carbon-containing compounds (e.g., organic solvent, organic residue on the catalyst surface, sacrificial agent, and carbon-based catalyst) in the reaction system have the possibility to be decomposed into small

fragments (e.g., CO, CH₄, and C₂H₄), and contribute to false-positive results. Therefore, systematic experiments should be carried out to identify the carbon source of the detected products, including i) control experiments in N₂ or Ar, ii) dark experiments at a similar temperature to photo-reaction, and iii) isotope experiments using $^{13}\text{CO}_2$ as the precursor. In addition, when pure water is employed as the reductant, oxidation products (i.e., O₂ and H₂O₂) should be examined to judge if the CO₂ conversion is realized by a catalytic reaction or a surface chemical reaction (i.e., OV is consumed and cannot be regenerated during the reaction). In detail, product balance should be reached for CO₂ reduction and H₂O oxidation (e.g., CO₂ + H₂O \rightarrow CO + 1/2 O₂ + H₂O). Only with all the above experiments being carefully implemented, it can be safe to conclude that the detected products are indeed from the photocatalytic reduction of CO₂.

Thirdly, the development of advanced CO₂ reduction systems to make efficient utilization of both photogenerated electrons and holes is pressing for achieving high overall solar-to-fuel conversion efficiency.

^b B: band structure, P: plasmonic effect.

^c M: mono-component metal oxide, H: hybrid nanostructure.

 $^{^{\}rm d}$ A: $\rm CO_2$ adsorption and activation, D: OV-based active site design.

^e Performance enhancement brought about by OV.

^f DMC: dimethyl carbonate.

⁸ The high reaction temperature is attributed to the photothermaleffect of the catalyst.

h TEOA: triethanolamine.

Specifically, the holes-related oxidative half-reaction should be well engineered given the slow kinetics of water oxidation. Currently, sacrificial agents (e.g., triethanolamine, methanol, and isopropanol) are commonly employed to accelerate the extraction of holes. However, this will lead to high capital costs and thus limit the practical application of photocatalytic CO₂ reduction. As a promising alternative, the combination of CO₂ reduction with selective organic oxidation can give rise to CO₂ reduction products and value-added chemicals simultaneously, achieving the double-win of catalytic performance and economic efficiency. In addition, the coupling of photocatalytic CO₂ reduction with plastic oxidation (i.e., reduction of the CO₂ obtained from plastic oxidation) has recently emerged as a feasible pathway for plastic upcycling. The knowledge obtained in these systems will also in turn contribute to the advancement of photocatalytic CO₂ reduction with H₂O.

Fourthly, the current methods for OV engineering are mainly focused on tuning the concentration and location (e.g., surface and bulk) of OV in mono-component metal oxide photocatalysts. The controllable formation of OV in more complicated systems (e.g., metal oxide-based hybrid nanostructures) remains a challenge, resulting in substantially limited mechanistic understanding and photocatalytic performance for CO2 reduction. In such a system, OV can be located at the material surface, bulk, and interface. Moreover, additional active sites can exist in the system to work together with OV for CO2 adsorption, activation, and conversion. In this context, developing controllable synthetic methods to maximize the synergistic effect between various OVs in the material and the additional active sites on the material surface is promising for obtaining high-performance photocatalysts. The precise integration of OV is also essential for obtaining more profound insights into the structure-performance relationship of photocatalytic CO₂ reduction over OV-containing metal oxide-based hybrid nanocatalysts.

Fifthly, advanced characterization techniques for studying the dynamic changes of OV during photocatalytic CO2 reduction are still inadequate. For the typical in situ characterization techniques (e.g., EPR, XPS, STM, DRIFTS) in photocatalytic CO_2 reduction, there is a huge disparity between the characterization and the real reaction conditions. For instance, the current in situ EPR characterization is mainly carried out in the gas-solid phase to investigate the influence of light irradiation on the electronic structure of OV. Likewise, in situ XPS can only be carried out at near-ambient pressure (typically <0.5 mbar). In addition, the water content during in situ DRIFTS measurements needs to be well controlled since IR light is highly sensitive to water. These limitations significantly interfere with our understanding of the actual role of OV in photocatalytic reactions. Thus, the development of novel reaction systems that mimic the current in situ characterization conditions is highly desired. The combination of multiple characterization techniques is also recommended for gaining a comprehensive understanding of the role of OV in photocatalytic CO2 reduction.

Nonetheless, these challenges provide plenty of opportunities for future work in this blooming research field. At the current pace of development and rise in the understanding of the role of OV, it is anticipated that the rational design and controllable synthesis of OV-containing metal oxide will contribute to the next generation of photocatalysts for efficient CO₂ reduction. Meanwhile, it is believed that more latent functions of OV will start to emerge with continuous exploration. We hope this review can inspire more creative work in the field of OV engineering for efficient photocatalytic CO₂ reduction.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

Data will be made available on request.

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